

Numerical analysis of the influence of solid-state phase transformations on the mechanical behavior of the Ti-6AI-4V alloy

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Selective Laser Melting

- Selective Laser Melting (SLM) is one of the main additive manufacturing processes for the production of metallic components
- It is characterized by extremely localized heat input and high temperature gradients
- The Ti-6AI-4V alloy is one of the most common materials processed by SLM



C. Galy, E. Le Guen, E. Lacoste, and C. Arvieu, "Main defects observed in aluminum alloy parts produced by SLM: From causes to consequences," Addit. Manuf., vol. 22, no. July 2017, pp. 165–175, 2018

Introduction

Microstructure of Ti-6Al-4V alloy

- Microstructure evolution can be quite complex
- Tipically, it contains α, β and α'
- Less frequently, α" martensite can be found





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Thermo-metallurgical-mechanical modelling



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M. Masoomi, S. M. Thompson, and N. Shamsaei, "Laser powder bed fusion of Ti-6AI-4V parts: Thermal modeling and mechanical implications," Int. J. Mach. Tools Manuf., vol. 118–119, no. April, pp. 73–90, 2017

Adapted Koistinen-Marburger equation

 $\beta \rightarrow \alpha' \begin{bmatrix} f_{\alpha'}(T + \Delta T) = \frac{f_{\alpha'}(T) - \gamma \Delta T \left(f_{\beta}(T_0) - f_{\beta_r} + f_{\alpha'}(T_0) \right)}{1 - \gamma \Delta T} & |\dot{T}| \ge 410 \text{ °C/s}, \ T \le M_s \end{bmatrix}$ $f_{\beta_r} = \begin{cases} f_{\beta}(T_0) & f_{\beta}(T_0) < 0.25 \\ 0.25 (1 - f_{\beta}(T_0)) & f_{\beta}(T_0) \ge 0.25 \end{cases} \begin{bmatrix} M_s = 650 \text{ °C} \\ \gamma = 0.015 \text{ °C}^{-1} \end{bmatrix}$



Adapted Johnson-Mehl-Avrami equation

 $f_{\beta}^{eq}(T) = f_{\beta,0}^{eq}(T)(1-f_{\alpha'})$

450

 $T[^{\circ}C]$

600

750

300

0

0

150

$$\begin{array}{c} \alpha' \to \alpha + \beta \\ \hline f_{\alpha'}(T + \Delta T) = 1 - \left[1 - \exp\left(-k_{1}\left(\zeta_{\alpha'} + \Delta t\right)^{n_{1}}\right)\right]\left(1 - f_{\alpha'}^{eq}\right) \\ \hline f_{\alpha'}(T + \Delta T) = -\Delta f_{\alpha'}(T + \Delta T) f_{\beta,0}^{eq}(T + \Delta T) \\ \hline \Delta f_{\beta}(T + \Delta T) = -\Delta f_{\alpha'}(T + \Delta T) f_{\beta,0}^{eq}(T + \Delta T) \\ \hline f_{\beta}(T + \Delta T) = \left[1 - \exp\left(-k_{2}\left(\zeta_{\beta}^{h} + \Delta t\right)^{n_{2}}\right)\right]f_{\beta}^{eq} \\ \hline f_{\beta}(T + \Delta T) = 1 - f_{\alpha'} - \left[1 - \exp\left(-k_{2}\left(\zeta_{\beta}^{c} + \Delta t\right)^{n_{2}}\right)\right]\left(1 - f_{\alpha'} - f_{\beta}^{eq}\right) \\ \hline f_{\beta}(T + \Delta T) = 1 - f_{\alpha'} - \left[1 - \exp\left(-k_{2}\left(\zeta_{\beta}^{c} + \Delta t\right)^{n_{2}}\right)\right]\left(1 - f_{\alpha'} - f_{\beta}^{eq}\right) \\ \hline f_{\beta}(T + \Delta T) = 1 - f_{\alpha'} - \left[1 - \exp\left(-k_{2}\left(\zeta_{\beta}^{c} + \Delta t\right)^{n_{2}}\right)\right]\left(1 - f_{\alpha'} - f_{\beta}^{eq}\right) \\ \hline f_{\beta}(T + \Delta T) = 1 - f_{\alpha'} - \left[1 - \exp\left(-k_{2}\left(\zeta_{\beta}^{c} + \Delta t\right)^{n_{2}}\right)\right]\left(1 - f_{\alpha'} - f_{\beta}^{eq}\right) \\ \hline f_{\beta}(T + \Delta T) = 1 - f_{\alpha'} - \left[1 - \exp\left(-k_{2}\left(\zeta_{\beta}^{c} + \Delta t\right)^{n_{2}}\right)\right]\left(1 - f_{\alpha'} - f_{\beta}^{eq}\right) \\ \hline f_{\beta}(T + \Delta T) = 1 - f_{\alpha'} - \left[1 - \exp\left(-k_{2}\left(\zeta_{\beta}^{c} + \Delta t\right)^{n_{2}}\right)\right]\left(1 - f_{\alpha'} - f_{\beta}^{eq}\right) \\ \hline f_{\beta}(T + \Delta T) = 1 - f_{\alpha'} - \left[1 - \exp\left(-k_{2}\left(\zeta_{\beta}^{c} + \Delta t\right)^{n_{2}}\right)\right]\left(1 - f_{\alpha'} - f_{\beta}^{eq}\right) \\ \hline f_{\beta}(T + \Delta T) = 1 - f_{\alpha'} - \left[1 - \exp\left(-k_{2}\left(\zeta_{\beta}^{c} + \Delta t\right)^{n_{2}}\right)\right]\left(1 - f_{\alpha'} - f_{\beta}^{eq}\right) \\ \hline f_{\beta}(T + \Delta T) = 1 - f_{\alpha'} - \left[1 - \exp\left(-k_{2}\left(\zeta_{\beta}^{c} + \Delta t\right)^{n_{2}}\right)\right]\left(1 - f_{\alpha'} - f_{\beta}^{eq}\right) \\ \hline f_{\beta}(T + \Delta T) = 1 - f_{\alpha'} - \left[1 - \exp\left(-k_{2}\left(\zeta_{\beta}^{c} + \Delta t\right)^{n_{2}}\right)\right]\left(1 - f_{\alpha'} - f_{\beta}^{eq}\right) \\ \hline f_{\beta}(T + \Delta T) = 1 - f_{\alpha'} - \left[1 - \exp\left(-k_{2}\left(\zeta_{\beta}^{c} + \Delta t\right)^{n_{2}}\right)\right]\left(1 - f_{\alpha'} - f_{\beta}^{eq}\right) \\ \hline f_{\beta}(T + \Delta T) = 1 - f_{\alpha'} - \left[1 - \exp\left(-k_{2}\left(\zeta_{\beta}^{c} + \Delta t\right)^{n_{2}}\right)\right]\left(1 - f_{\alpha'} - f_{\beta}^{eq}\right) \\ \hline f_{\beta}(T + \Delta T) = 1 - f_{\alpha'} - \left[1 - \exp\left(-k_{2}\left(\zeta_{\beta}^{c} + \Delta t\right)^{n_{2}}\right)\right]\left(1 - f_{\alpha'} - f_{\beta}^{eq}\right) \\ \hline f_{\beta}(T + \Delta T) = 1 - f_{\alpha'} - \left[1 - \exp\left(-k_{2}\left(\zeta_{\beta}^{c} + \Delta t\right)^{n_{2}}\right)\right]\left(1 - f_{\alpha'} - f_{\beta}^{eq}\right) \\ \hline f_{\beta}(T + \Delta T) = 1 - f_{\alpha'} - \left[1 - \exp\left(-k_{2}\left(\zeta_{\beta}^{c} + \Delta t\right)^{n_{2}}\right)\right]\left(1 - f_{\alpha'} - f_{\beta}^{eq}\right) \\ \hline f_{\beta}(T + \Delta T) = 1 - \left[1 - \exp\left(-k_{2}\left(\zeta_{\beta}^{c} + \Delta t\right)^{n_{2}}\right)\right]\left(1 - f_{\alpha'} - f_{\beta}^{eq}\right) \\ \hline f_{\beta}(T +$$

900

6

Thermal and solid-state transformation induced strains

$$\boldsymbol{\varepsilon} = \boldsymbol{\varepsilon}_{e} + \boldsymbol{\varepsilon}_{p} + \boldsymbol{\varepsilon}_{th} + \boldsymbol{\varepsilon}_{sst}$$

• Since the lattice parameters of the α and α' phases are similar, $\alpha_{\alpha} = \alpha_{\alpha'}$

$$\Delta \boldsymbol{\varepsilon}^{\text{th}} = \begin{cases} \begin{bmatrix} \alpha_{\beta} f_{\beta} + \alpha_{\alpha} (1 - f_{\beta}) \end{bmatrix} \Delta T \mathbf{I} & \text{No transformation} \\ \begin{bmatrix} \alpha_{\beta} f_{\beta} + \alpha_{\alpha} (1 - f_{\beta} - \Delta f_{\alpha'}) \end{bmatrix} \Delta T \mathbf{I} & \beta \rightarrow \alpha' \\ \begin{bmatrix} \alpha_{\beta} f_{\beta} + \alpha_{\alpha} (1 - f_{\beta} - \Delta f_{\alpha}) \end{bmatrix} \Delta T \mathbf{I} & \beta \rightarrow \alpha \\ \begin{bmatrix} \alpha_{\beta} (f_{\beta} - \Delta f_{\beta}) + \alpha_{\alpha} (1 - f_{\beta} - \Delta f_{\alpha,m}) \end{bmatrix} \Delta T \mathbf{I} & \alpha \rightarrow \beta \text{ or } \alpha' \rightarrow \alpha + \beta \end{cases}$$
$$\Delta \boldsymbol{\varepsilon}^{\text{sspt}} = \boldsymbol{\varepsilon}^{\Delta V}(T) \Delta f_{i}(T) \mathbf{I}$$



J. W. Elmer, T. A. Palmer, S. S. Babu, and E. D. Specht, "In situ observations of lattice expansion and transformation rates of α and β phases in Ti-6AI-4V," Mater. Sci. Eng. A, vol. 391, no. 1–2, pp. 104–113, 2005

Numerical model

- Single finite element
- Two sequential heating/cooling cycles replicate the laser movement
- Material is initially in powder state



Material properties

Temperature independent

Temperature dependent



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Phase volume fractions



Numerical results

Volume change relative to the initial volume



Simulation 1

• C1: expansion during cooling due to $\beta \rightarrow \alpha'$

Simulation 2

• C2: expansion during cooling due to $\beta \rightarrow \alpha'$

Numerical results

Stress evolution



- Simulation 6 (no SST) drastically overestimates the stress
- At the end of simulations 4 and 5, the differences in phase volume fractions yield a stress difference of nearly 92 MPa

- The prediction of solid-state phase transformations is important for an accurate estimation of the material's volume change and stress field
- In the simulations that account for solid-state phase transformations, changing the heating/cooling rates yielded completely different final solid phase volume fractions
- The predicted volume change of the material relative to its initial volume showed a 0.5% difference
- In terms of the final stress value in the restricted model, this difference corresponds to a stress discrepancy of 92 MPa
- When solid-state phase transformations were not considered, the predicted stress was drastically overestimated (> 1000 MPa)

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Thank you for your attention!